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DISPERSAL OF MOLECULAR CLOUDS BY PHOTOIONISATION AND RADIATION PRESSURE

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Slow and Inefficient SF

• SF is slow: $t_{
m dep} = M_{
m cl}/
m SFR$ $\gg t_{
m ff} \sim 1-10~
m Myr$

(Bigiel+08, Kennicutt+12, Leroy+13)

• SF is inefficient:

$$SFE = \frac{M_*}{M_* + M_{cl}}$$



- Star-forming clouds in the Milky Way have SFE ~ 0.002-0.2, a few % on average (Myers+86, Williams+97, Carpenter00, Murray+11, Garcia+14)
- Shorter t_{dep}, higher SFE in dese, high-Σ environments (Meier+02, Turner+15; Leroy+13, 15)

Lifetime and Disruption of GMCs



No **Massive SF** only **H II regions** H II regions + young clusters lacksquare

only

young clusters

- Final (or net) SFE for steady SFR SFE_{final} = $\frac{SFR \times t_{cl}}{M_{cl}} = \frac{t_{cl}}{t_{dep}}$
- H II regions and young star clusters seen with GMCs
- Each type represents evolutionary sequence
- Lifetime of GMCs: 20-30 Myr

(see also Murray+11, Miura+12, Meidt+15)

Blitz+07; Kawamura+09



Radiation feedback can remove residual gas around young star clusters, controlling SFE and lifetime of molecular clouds.



Destructive Effects of UV Radiation



- Photoionization
 - Photoevaporation
 - Thermal pressure
 - Rocket effect
 - Radiation pressure on dust
 - Important in massive, high-Σ clouds

(Whitworth+79, Franco+94, Williams+97, Matzner+02, Krumholz+06, Krumholz+09, Murray+10, Goldbaum+11, Lopez+11, Zamora-Aviles+12, Kim+16)

Previous Models on Cloud Disruption

Photoionization only (Dale+12, 13)



(see also Walch+12; Geen+15, 16, 17; Howard+16, 17; Gavagnin+17)

Previous Models on Cloud Disruption

Radiation pressure only (Raskutti+16)



- SFE_{final} of radiation pressure-regulated clouds depend on (the distribution of) $\Sigma_{\rm cl}$

(Fall+10, Thompson+16, Kim+16, Raskutti+16, Grudic+17)

Key Issues

- Can UV radiation feedback sustain cloud turbulence or completely destroy cloud?
- Can UV radiation feedback explain low SFE of GMCs?
- What is the timescale for gas dispersal?
- Relative importance of mass loss mechanisms
 - Photoevaporation
 - Dynamical ejection
- Escape fraction of ionizing radiation
- Boundedness of star clusters

3D radiation hydrodynamic simulations of star cluster formation in turbulent clouds including both photoionization and radiation pressure

Numerical Method







- Grid-based code Athena (Stone+08)
- Uniform density sphere with initial injection of turbulence (Stone+98)
- Marginally bound with $\,\alpha_{\rm vir}=2$
- Star formation and accretion via sink particle method (Gong+13)
- Temperature as a function of H-ionization fraction (20 K < T < 8000 K)

Radiative Transfer for Multiple Moving Point Sources



- Sink particles representing subclusters emit ionizing and non-ionizing radiation.
- Adaptive ray tracing (Abel+02) with improved parallel performance (Rosen+17, Kim+17, submitted)

Cloud Parameters





Emission measure (Simulation)



NGC 602 (HST image)

 Irregular structures such as "fingers", "elephant trunks" naturally arise from turbulent density field sculpted by UV radiation (e.g., Mellema+06, Arthur+11, Tremblin+12, Gritschneder+09, 10)

Evolution of Key Characteristics



- SFE_{final}: 15%
- Photoevaporation is the dominant mass loss mechanism $(M_{\rm ev}/M_{\rm cl} \sim 0.8)$

Final SFE and Evaporation Efficiency



 SFE_{final} and evaporation efficiency depend primarily on the initial cloud surface density

Characterizing Mass Loss by Photoevaporation



Ionization-recombination balance

$$Q_{\rm eff} \approx \alpha_{\rm i} n_{\rm i}^2 A H$$

- Q_{eff} : ionizing photon absorption rate A : area of ionization front
- H : thickness of recombination layer
- Mass loss rate

$$\dot{M}_{\rm ev} = (\mu n_{\rm i} c_{\rm i}) \times A$$
$$= \mu c_{\rm i} \left(\frac{Q_{\rm eff}}{\alpha_{\rm i} A H}\right)^{1/2} \times A$$
$$f \sim Q$$
$$\sim R_{\rm cl} \qquad \propto (Q_{\rm eff} A/H)^{1/2}$$
$$\sim 4\pi R_{\rm cl}^2 \qquad \propto Q^{1/2} R_{\rm cl}^{1/2}$$

Mass Loss Rate and Evaporation Time Scale



- Mass loss rate $\propto Q_i^{1/2} R_{cl}^{1/2}$
- Evaporation timescale (measured in units of free-fall time) depends primarily on Σ_{cl}

Control of SFE by Photoevaporation

• Stellar mass required to photoevaporate $M_{ev} = M_{cl} - M_*$

$$M_{\rm cl} = M_* + M_{\rm ev} = M_* + \langle \dot{M}_{\rm ev} \rangle t_{\rm ev} = M_* + \frac{\mu_{\rm H} c_{\rm i}}{\alpha_{\rm i}^{1/2}} \left\langle \frac{Q_{\rm eff}^{1/2} A^{1/2}}{H^{1/2}} \right\rangle t_{\rm ev}$$

• Dividing by M_{cl}

$$1 = \varepsilon_* + \phi \left(\frac{t_{\rm ev}}{t_{\rm ff,0}}\right) \left(\frac{\Sigma_{\rm cl}}{500 \ M_{\odot} {\rm pc}^{-2}}\right)^{-1} \varepsilon_*^{1/2}$$
$$\int_{\sim 0.3}^{1/2} \phi_t = -1.65 + 1.55 \log_{10} \Sigma_{\rm cl,0}$$

Final SFE and Evaporation Efficiency



-10

-10

0

x [pc]

10

 Rocket effect becomes important for low-mass clouds with

$$M_{\rm ej} v_{\rm esc} \propto M_{\rm ev} c_{\rm i}$$

Discussion

- t_{ev} is a good measure of the timescale for gas dispersal after the onset of massive SF for photoevaporationdominated cases
 - t_{ev} ranges between ~1-3 t_{ff} or 2-10 Myr
 - $t_{dep} \sim t_{cl}/SFE_{final} \sim t_{ev}/SFE_{final} \sim 10-100 \text{ Myr}$
- Possible reasons for higher SFE and shorter t_{dep} than observed?
 - Artificial initial condition
 - Neglect of stellar winds and supernovae
 - Absence of magnetic support
 - Unresolved, subgrid-scale physical processes
 - Need for a subgrid model for sink particles? (Howard+14,16)

Mass Loss Rate of Evaporating Pillars

$$\dot{M}_{\rm ev,obs} = \pi r^2 m_{\rm H} n_{\rm H} v$$
$$\dot{M}_{\rm ev,theory} \propto \Phi^{1/2} A r^{-1/2}$$
$$\propto \Phi^{1/2} r^{3/2}$$

(e.g., Kahn+69, Bertoldi89,+90, Lefloch+94)





Analogy between Pillars and Clouds



Summary

- Radiation hydrodynamic simulations of star cluster formation in turbulent GMCs
 - 1. Photoevaporation plays a dominant role in the disruption of GMCs typical of the MW.
 - 2. High SFE_{final} for high-Σ clouds
 - Cloud destruction occurs within a few dynamical timescales once sufficiently luminous H II regions form.